

CHAPTER III

DYNAMICAL INHERITANCE PROPERTY
AND
MAGNETO-DUST DISTRIBUTION.

1. INTRODUCTION :

The critical study of the work due to Maarteen and Maharaj (1986) suggests that there do exist proper Conformal Killing vectors in FRW models and perfect fluid distribution. This suggested clue to Duggal (1992) to modify the concept of curvature collineation suitable for proper Conformal Killing vector and other related proper symmetries. Accordingly he introduced a new symmetry called as Curvature Inheritance defined as,

$$\mathfrak{L}_{\xi} R^a_{bcd} = 2 \alpha R^a_{bcd}.$$

This implies Ricci Inheritance property,

$$\mathfrak{L}_{\xi} R_{ab} = 2 \alpha R_{ab}.$$

The prime goal of this work is to explore the geometrical and dynamical properties of Dynamical Inheritance associated with Magneto-Dust Distribution described by energy momentum tensor I (3.4) through the definition given below :

2. DEFINITION : DYNAMICAL INHERITANCE :

The space-time is said to admit Dynamical Inheritance along vector field ξ if,

$$\mathfrak{L}_{\xi} T_{ab} = 2 \alpha T_{ab} . \quad 1$$

where α is scalar function of co-ordinates and T_{ab} is the stress energy tensor.

This stress energy tensor I (3.4) involves two eigen vectors one of which is the time-like eigen vector u^a with eigen value $e_1 = \rho + \frac{1}{2} \mu h^2$ and the other is the space-like eigen vector h^a with eigen value

$$e_2 = -\frac{1}{2} \mu h^2 = \frac{1}{2} \mu h^2$$

Hence the study of Dynamical Inheritance described by (1) is governed by the choice of the Inheritance vector ξ with two separate cases as

$$\xi^a = u^a \quad \text{and} \quad \xi^a = h^a$$

Equation (1) can be expressed as

$$(T_{ab;c}) \xi^c + (T_{db})(\xi^d_{;a}) + (T_{ad})(\xi^d_{;b}) = 2 \alpha T_{ab} \quad 2$$

We study this expression with the following two specific types of vector ξ

CASE I : If $\xi^a = u^a$

For this choice equation (2) can be written as

$$(T_{ab;c}) u^c + (T_{db})(u^d_{;a}) + (T_{ad})(u^d_{;b}) = 2 \alpha T_{ab} \quad 3$$

Theorem 1 : If the Magneto-Dust Distribution admits Dynamical Inheritance property along flow lines then the active gravitational mass density is inherited iff $(\mu h^2)' + \dot{\mu} h^2 + \mu h^2 \theta = 0$.

Proof : If equation (3) is transvected with g^{ab} , we get

$$\begin{aligned} (g^{ab} T_{ab;c}) u^c + (T_{db} g^{ab})(u^d_{;a}) + (g^{ab} T_{ad})(u^d_{;b}) &= \\ &= 2 \alpha g^{ab} T_{ab} \quad \because g^{ab}_{;c} = 0 \end{aligned}$$

$$\text{i.e. } (T_{;c}) u^c + T^a_d (u^d_{;a}) + T^b_d (u^d_{;b}) = 2 \alpha T,$$

$$\text{i.e. } (T_{;c}) u^c + 2 T^a_d (u^d_{;a}) = 2 \alpha T,$$

$$\text{i.e. } (T_{;c}) u^c + 2 T^{ad} (u_{d;a}) = 2 \alpha T. \quad 4$$

Putting $T = \rho$ (vide I - 4.7) and using expression of T^{ad} (vide I - 3.4) in equation (4), we have

$$\begin{aligned} (\rho_{;c}) u^c + 2 [(\rho + \mu h^2) u^a u^d - \frac{1}{2} \mu h^2 g^{ad} - \mu h^a h^d] \\ [u_{d;a}] = 2 \alpha \rho. \end{aligned} \quad 5$$

We have $u^a_{;a} = \theta$, $\rho_{;c} u^c = \dot{\rho}$, $u^d_{;a} u^a_d = 0$.

Hence equation (5) can be written as

$$\dot{\rho} + 2 [(\rho + \mu h^2) u^a u^d u^a_{;d} - \frac{1}{2} \mu h^2 g^{ad} u^a_{;d} - \mu h^a h^d u^a_{;d}] = 2 \alpha \rho,$$

$$\text{i.e. } \dot{\rho} - \mu h^2 \theta - 2 \mu u_d{}_{;a} h^d h^a = 2 \alpha \rho. \quad 6$$

Maxwell equation I (8.7) direct

$$\mu u^a_{;d} h^d h^a = -\dot{\mu} h^2 - \mu h^2 \theta - (\mu/2) (h^2)' \quad 7$$

Hence from (7), equation (6) can be written as

$$\dot{\rho} - \mu h^2 \theta - 2 [-\dot{\mu} h^2 - \mu h^2 \theta - (\mu/2) (h^2)'] = 2 \alpha \rho, \quad 8$$

$$\text{i.e. } \dot{\rho} + \mu h^2 \theta + 2 \dot{\mu} h^2 + \mu (h^2)' = 2 \alpha \rho, \quad 9$$

$$\text{i.e. } [\rho + \mu h^2]' + \dot{\mu} h^2 + \mu h^2 \theta = 2 \alpha \rho, \quad 10$$

$$\text{i.e. } L_u (\rho + \mu h^2) = 2 \alpha \rho - \dot{\mu} h^2 - \mu h^2 \theta, \quad 11$$

This implies that

$$L_u \rho = 2 \alpha \rho \quad \text{iff} \quad (\mu h^2)' + \dot{\mu} h^2 + \mu h^2 \theta = 0, \quad 12$$

But $\rho = T$ vide I (4.7),

Hence from (12), we get

$$L_u T = 2 \alpha T \quad \text{iff} \quad (\mu h^2)' + \dot{\mu} h^2 + \mu h^2 \theta = 0.$$

Hence the theorem. $\square$

Theorem 2 : If the Magneto-Dust Distribution admits Dynamical Inheritance property along flow lines then time-like eigen value is inherited along flow lines.

Proof : On multiplying equation (3) by $u^a u^b$, we get

$$\begin{aligned} u^a u^b (T_{ab;c}) u^c + u^a u^b (u^d_{;a}) T_{ab} + u^a u^b (u^d_{;b}) T_{ab} \\ = 2 \alpha u^a u^b T_{ab} \end{aligned} \quad 13$$

The three terms on left hand side of equation (13) are simplified in following manner.

$$\begin{aligned} \text{Consider, } u^a u^b (T_{ab;c}) u^c &= \\ &= u^a u^b [(\rho + \mu h^2) u_a u_b - \frac{1}{2} \mu h^2 g_{ab} - \mu h_a h_b] ;_c u^c \\ &\quad \text{vide I (3.4)} \\ &= u^a u^b [(\rho + \mu h^2) u_a u_b] ;_c u^c - (\frac{1}{2} \mu h^2) ;_c u^a u^b g_{ab} u^c \\ &\quad - (\mu h_a h_b) ;_c u^c u^a u^b u^c, \end{aligned} \quad 14$$

Using $u_a ;_c u^a = 0$, $u^a h_a = 0$, $g^{ab} ;_c = 0$ and $u^a u_a = 1$ in equation (14), we get

$$\begin{aligned} u^a u^b (T_{ab;c}) u^c &= \\ &= (\rho + \mu h^2) ;_c u^c + (\rho + \mu h^2) u_a u_b u_a ;_c u_b u^c + \\ &\quad + (\rho + \mu h^2) u^a u^b u_b ;_c u_a u^c - \\ &\quad - \frac{1}{2} [g_{ab} (\mu ;_c u^a u^b h^2 + h^2 ;_c u^a u^b \mu)] u^c \\ &\quad - [\mu ;_c u^a u^b h_a h_b + \mu h_a ;_c u^a u^b h_b + \mu h_b ;_c h_a u^a u^b] u^c, \\ &= (\rho + \mu h^2) ;_c u^c - \frac{1}{2} (\mu ;_c h^2 + h^2 ;_c \mu) u^c \\ &= (\rho + \mu h^2) ;_c u^c - \frac{1}{2} (\mu h^2) ;_c u^c. \end{aligned}$$

Finally we get

$$u^a u^b (T_{ab;c}) u^c = (\rho + \frac{1}{2} \mu h^2) ;_c u^c. \quad 15$$

Now consider

$$\begin{aligned} u^a u^b (u^d_{;a}) T_{ab} &= \\ &= u^a u^b [(\rho + \mu h^2) u_d u_b - \frac{1}{2} \mu h^2 g_{ab} - \mu h_d h_b] u^d_{;a} \\ &\quad \text{vide I (3.4)} \\ &= (\rho + \mu h^2) u^a u^b u_d u_b u^d_{;a} - \frac{1}{2} \mu h^2 g_{ab} u^a u^b u^d_{;a} - \\ &\quad - \mu h_d h_b u^a u^b u^d_{;a}, \end{aligned} \quad 16$$

We have $u^d{}_{;a}u_d = 0$, $u^a u_a = 1$, $u^b h_b = 0$,

Hence from equation (16), we get

$$u^a u^b (u^d{}_{;a}) T_{db} = 0. \quad 17$$

Further consider

$$\begin{aligned} u^a u^b (u^d{}_{;b}) T_{ad} &= u^a u^b [(\rho + \mu h^2) u_a u_d \\ &\quad - \frac{1}{2} \mu h^2 g_{ad} - \mu h_a h_d] u^d{}_{;b} \\ &\quad \text{vide I (3.4)} \\ &= (\rho + \mu h^2) u^a u^b u_a u_d u^d{}_{;b} \\ &\quad - \frac{1}{2} \mu h^2 u^a u^b g_{ad} u^d{}_{;b} - \\ &\quad - \mu h_a h_d u^a u^b u^d{}_{;b}, \\ u^a u^b (u^d{}_{;a}) T_{ad} &= 0. \quad 18 \\ \therefore u^d{}_{;b} u_d &= 0, u^a u_a = 1, u^a h_a = 0 \end{aligned}$$

Thus from (15), (17) and (18) equation (13) reduces as

$$(\rho + \frac{1}{2} \mu h^2)_{;c} u^c = 2 \alpha u^a u^b T_{ab}, \quad 19$$

This implies that

$$(\rho + \frac{1}{2} \mu h^2)_{;c} u^c = 2 \alpha (\rho + \frac{1}{2} \mu h^2) \quad \text{vide I (4.2)}$$

$$\text{i.e. } \frac{L}{u} (\rho + \frac{1}{2} \mu h^2) = 2 \alpha (\rho + \frac{1}{2} \mu h^2), \quad 20$$

$$\text{i.e. } \frac{L}{u} e_1 = 2 \alpha e_1, \quad \text{vide I (4.2)} \quad 21$$

Here the proof of theorem is completed.

Note 1 : Equation (20) can be rearranged as

$$\dot{\rho} + \frac{1}{2} \dot{\mu} h^2 + \frac{1}{2} \mu (h^2)' = 2 \alpha (\rho + \frac{1}{2} \mu h^2). \quad 22$$

Making use of continuity equation I (9.10) in (22), we get

$$\begin{aligned} (\frac{1}{2} \dot{\mu} h^2 - \rho \theta) + \frac{1}{2} \dot{\mu} h^2 + \frac{1}{2} \mu (h^2)' &= \\ &= 2 \alpha (\rho + \frac{1}{2} \mu h^2). \end{aligned}$$

$$\text{i.e. } \dot{\mu} h^2 - \rho \theta + \frac{1}{2} \mu (h^2)' = 2 \alpha (\rho + \frac{1}{2} \mu h^2). \quad 23$$

This equation (23) exhibits the effect of dynamical entities on Inheritance scalar function α .

Theorem 3 : If the Magneto-Dust Distribution admits Dynamical Inheritance property along flow lines then

$$\underset{u}{L} \mu h^2 = \alpha \mu h^2 \quad \text{iff} \quad \mu h^2 + 2 \mu h^2 \theta = 0.$$

Proof : On transvecting equation (3) with $h^a h^b$, we get

$$\begin{aligned} h^a h^b (T_{ab;c}) u^c + h^a h^b T_{db} (u^d{}_{;a}) + h^a h^b T_{ad} (u^d{}_{;b}) \\ = 2 \alpha h^a h^b T_{ab} \end{aligned} \quad 24$$

Each term on left hand side of equation (24) is

simplified in following manner.

$$\begin{aligned} \text{Consider } h^a h^b (T_{ab;c}) u^c &= \\ &= h^a h^b [(\rho + \mu h^2) u_a u_b - \frac{1}{2} \mu h^2 g_{ab} - \mu h_a h_b]; u^c \\ &\quad \text{vide I (3.4)} \\ &= [(\rho + \mu h^2);_c u^c u_a u_b h^a h^b + (\rho + \mu h^2) u_a;_c u_b h^a h^b \\ &\quad + (\rho + \mu h^2) u_b;_c u_a h^a h^b] u^c - \frac{1}{2} [\mu h^2;_c h^a h^b + \mu;_c h^2 h^a h^b] \\ &\quad g_{ab} u^c - [\mu;_c h_a h_b h^a h^b + \mu h_a;_c h_b h^a h^b + \mu h_b;_c h_a h^a h^b] u^c. \end{aligned} \quad 25$$

We have $u_a;_c u^a = 0$, $h_a;_c u^c = h_a$, $h_a h^a = -\frac{1}{2} (h^2)$, $u^a h_a = 0$, $h^a h_a = -h^2$,

$\rho;_c u^c = \dot{\rho}$ Hence from equation (25)

$$\begin{aligned} h^a h^b (T_{ab;c}) u^c &= \\ &= -\frac{1}{2} [\mu;_c u^c h^2 h^a h_a + \mu h^2;_c u^c h^a h_a] \end{aligned}$$

$$\begin{aligned}
& - [\mu_{;c} u^c h^4 + \mu h_a{}_{;c} h^a{}_{;c} h^a (-h^2) u^c \\
& + \mu h_b{}_{;c} u^c (-h^2) h^b] ,
\end{aligned} \tag{26}$$

$$\begin{aligned}
\text{i.e. } h^a h^b (T_{ab;c}) u^c &= \\
&= - \frac{1}{2} [\dot{\mu} h^2 (-h^2) + \mu (h^2)' (-h^2)] \\
&\quad - [\dot{\mu} h^4 + \mu (h_a^a) h^a (-h^2) + \mu (h_b^b) h^b (-h^2)]
\end{aligned} \tag{27}$$

$$\begin{aligned}
&= \frac{1}{2} \dot{\mu} h^4 + \frac{1}{2} \mu h^2 (h^2) - \mu h^4 - \frac{1}{2} \mu h^2 (h^2)' \\
&\quad - \frac{1}{2} \mu h^2 (h^2)' \\
&= - \frac{1}{2} \dot{\mu} h^4 - \frac{1}{2} \mu h^2 (h^2)' \\
&= - \frac{1}{2} h^2 [\dot{\mu} h^2 + \mu (h^2)'] \\
h^a h^b (T_{ab;c}) u^c &= - \frac{1}{2} h^2 [\mu h^2]' .
\end{aligned} \tag{28}$$

Now consider

$$\begin{aligned}
h^a h^b T_{db} (u^d{}_{;a}) &= \\
&= h^a h^b [(\rho + \mu h^2) u_d u_b - \frac{1}{2} \mu h^2 g_{db} - \mu h_d h_b] u^d{}_{;a} \\
&\quad \text{vide I (3.4)} \\
&= (\rho + \mu h^2) u_d u_b h^a h^b - \frac{1}{2} \mu h^2 g_{db} h^a h^b - \\
&\quad - \mu h^a h^b h_d h_b] u^d{}_{;a} , \\
&= (\rho + \mu h^2) u^d{}_{;a} u_d u_b h^a h^b - \frac{1}{2} \mu h^2 u^d{}_{;a} g_{db} h^a h^b - \\
&\quad - \mu u^d{}_{;a} h^a h^b h_d h_b .
\end{aligned} \tag{29}$$

We have $h_a{}_{;c} u^c = \dot{h}_a$, $\dot{h}_a h^a = -\frac{1}{2} (h^2)$, $u^a h_a = 0$, $h^a h_a = -h^2$, $\rho_{;c} u^c = \dot{\rho}$

Hence from equation (29), we get

$$\begin{aligned}
h^a h^b T_{db} (u^d{}_{;a}) &= - \frac{1}{2} \mu h^2 u^d{}_{;a} h^a h_d + \mu h^2 u^d{}_{;a} h^a h_d \\
&= \frac{1}{2} \mu h^2 [u_a{}_{;d} h^a h^d] \\
&= \frac{1}{2} h^2 [- \dot{\mu} h^2 - \mu h^2 \theta - (\mu/2) (h^2)']
\end{aligned} \tag{30}$$

vide I (8.7)

$$h^a h^b (T_{db}) u^d{}_{;a} = -h^2/4 [2 \dot{\mu} h^2 + 2 \mu h^2 \theta + \mu (h^2)'] \quad 31$$

Further consider,

$$h^a h^b T_{ad} (u^d{}_{;b}) = h^a h^b [(\rho + \mu h^2) u_a u_d - \frac{1}{2} \mu h^2 g_{ad} - \mu h_a h_d] u^d{}_{;b} \quad 32$$

$$= [(\rho + \mu h^2) u^d{}_{;b} h^a h^b u_a u_d - \frac{1}{2} \mu h^2 g_{ad} h^a h^b u^d{}_{;b} - \mu u^d{}_{;b} h^a h^b h_a h_d] \quad 33$$

If we put $u^a h_a = 0$, $h^a h_a = -h^2$, $h_a h^a = -\frac{1}{2} (h^2)'$, $u^d{}_{;b} u_d = 0$

in equation (33), we get

$$\begin{aligned} h^a h^b T_{ad} (u^d{}_{;b}) &= -\frac{1}{2} \mu h^2 u^d{}_{;b} h_d h^b + h^2 u^d{}_{;b} h^b h_d, \\ h^a h^b T_{ad} (u^d{}_{;b}) &= \frac{1}{2} \mu h^2 u_d{}_{;b} h^d h^b, \\ &= \frac{1}{2} h^2 [-\mu h^2 - \mu h^2 \theta - (\mu/2) (h^2)'], \\ &\quad \text{vide I (8.7)} \\ &= -h^2/4 [2 \dot{\mu} h^2 + 2 \mu h^2 \theta + \mu (h^2)'] \quad 34 \end{aligned}$$

Thus from (28), (31) and (34), equation (24) reduces to

$$\begin{aligned} &-\frac{1}{2} h^2 (\mu h^2)' - \frac{1}{2} h^2 [2 \dot{\mu} h^2 + 2 \mu h^2 \theta + \mu (h^2)'] \\ &= 2 \alpha h^a h^b T_{ab}, \quad 35 \end{aligned}$$

$$\begin{aligned} \text{i.e. } &-\frac{1}{2} h^2 (\mu h^2)' - \frac{1}{2} h^2 [2 \dot{\mu} h^2 + 2 \mu h^2 \theta + \mu (h^2)'] \\ &= -2 \alpha h^4, \quad \text{vide I (4.4)} \end{aligned}$$

$$2 (\mu h^2)' + \dot{\mu} h^2 + 2 \mu h^2 \theta = 2 \alpha \mu h^2 \quad 36$$

This implies that

$$2 L_u (\mu h^2) = 2 \alpha \mu h^2 - \dot{\mu} h^2 - 2 \mu h^2 \theta. \quad 37$$

Hence the proof of theorem is completed.

Theorem 4 : If Magneto-Dust Distribution admits Dynamical inheritance property along flow lines then acceleration is orthogonal to magnetic lines.

Proof : On contracting equation (2) with $u^a h^b$, we get

$$\begin{aligned} u^a h^b (T_{ab;c}) u^c + u^a h^b T_{db} (u^d;_a) + T_{ad} u^a h^b (u^d;_b) \\ = 2 \alpha u^a h^b T_{ab}, \end{aligned} \quad 38$$

$$\begin{aligned} \text{i.e. } u^a h^b (T_{ab;c}) u^c + T_{db} u^a h^b (u^d;_a) + T_{ad} u^a h^b (u^d;_b) = 0. \\ \text{vide I (4.5)} \end{aligned} \quad 39$$

The three terms on left hand side of (39) are simplified in following manner :

Consider,

$$\begin{aligned} u^a h^b (T_{ab;c}) u^c &= u^a h^b [(\rho + \mu h^2) u_a u_b \\ &\quad - \frac{1}{2} \mu h^2 g_{ab} - \mu h_a h_b] ;_c u^c, \\ &= [(\rho + \mu h^2) ;_c u^a h^b u_a u_b + (\rho + \mu h^2) u_a ;_c \\ &\quad u_b u^a h^b + (\rho + \mu h^2) u_b ;_c u_a u^a h^b] u^c \\ &\quad - \frac{1}{2} [g_{ab} \mu ;_c h^2 u^a h^b + g_{ab} \mu h^2 ;_c u^a h^b] u^c \\ &\quad - [\mu ;_c u_a h^b h_a h_b + \mu h_a ;_c h_b u^a h^b + \\ &\quad + \mu h_b ;_c h_a u^a h^b] u^c. \end{aligned} \quad 40$$

Using $u_a h^a = 0$, $\dot{u}_a = u_a ;_c u^c$, $\dot{u}_a u^a = 0$ in (40), we get

$$\begin{aligned} u^a h^b (T_{ab;c}) u^c &= (\rho + \mu h^2) \dot{u}_b h^b + \mu h^2 h_a \dot{u}^a, \\ &= (\rho + \mu h^2) \dot{u}_b h^b - \mu h^2 \dot{u}^a h_a, \\ &\quad \because \dot{u}_a h^a = -h_a \dot{u}^a \\ &= (\rho) \dot{u}_b h^b. \end{aligned} \quad 41$$

Now consider,

$$\begin{aligned} u^a h^b T_{db} u^d ;_a &= u^a h^b [(\rho + \mu h^2) u_d u_b \\ &\quad - \frac{1}{2} \mu h^2 g_{db} - \mu h_d h_b] u^d ;_a, \end{aligned}$$

$$\begin{aligned}
&= (\rho + \mu h^2) u^d{}_{;a} u^a h^b u_d u_b \\
&\quad - \frac{1}{2} \mu h^2 u^d{}_{;a} u^a h^b g_{db} - \mu u^d{}_{;a} u^a h^b h_d h_b
\end{aligned} \tag{42}$$

Using $u^a u_a = 1$, $u^a h_a = 0$, $u^d{}_{;a} u^a = \dot{u}^d$, $u^d{}_{;a} u_d = 0$ in (42), we get

$$\begin{aligned}
u^a h^b T_{db} u^d{}_{;a} &= -\frac{1}{2} \mu h^2 u^d{}_{;a} u^a h_d - \mu \dot{u}^d (-h^2) h_b, \\
&= -\frac{1}{2} \mu h^2 \dot{u}^d h_d + \mu h^2 \dot{u}^d h_b, \\
&= +\frac{1}{2} \mu h^2 \dot{u}^b h_b.
\end{aligned} \tag{43}$$

Further consider,

$$\begin{aligned}
u^a h^b T_{ad} u^d{}_{;b} &= u^a h^b [(\rho + \mu h^2) u_a u_d \\
&\quad - \frac{1}{2} \mu h^2 g_{ad} - \mu h_a h_d] u^d{}_{;b}, \\
&= (\rho + \mu h^2) u^d{}_{;b} u^a h^b u_a u_d \\
&\quad - \frac{1}{2} \mu h^2 u^d{}_{;b} u^a h^b g_{ad} - \mu u^d{}_{;b} u^a h^b h_a h_d,
\end{aligned} \tag{44}$$

We have $u^a u_a = 1$, $u^a h_a = 0$, $u^a{}_{;b} u_a = 0$. Hence from (44), we get

$$u^a h^b T_{ad} u^d{}_{;b} = 0. \tag{45}$$

Hence from (41), (43), and (45), equation (39) yields as,

$$\begin{aligned}
\rho \dot{u}_b h^b + \frac{1}{2} \mu h^2 \dot{u}^b h_b &= 0. \\
\text{i.e. } \rho \dot{u}_b h^b + \frac{1}{2} \mu h^2 \dot{u}^b h_b &= 0. \\
\text{i.e. } (\rho + \frac{1}{2} \mu h^2) \dot{u}_b h^b &= 0.
\end{aligned}$$

This implies that

$$\dot{u}_b h^b = 0, \quad \because \rho + \frac{1}{2} \mu h^2 \neq 0 \tag{46}$$

This completes the proof.

CASE II : If $\xi^a = h^a$.

For this choice equation (2) can be written as

$$(T_{ab;c})h^c + (T_{db})(h^d;_a) + (T_{ad})(h^d;_b) = 2\alpha T_{ab}. \quad 47$$

Theorem 5 : If the Magneto-Dust Distribution admits Dynamical Inheritance property along magnetic lines then the time eigen value is inherited along magnetic lines iff $h^d u_d = 0$.

Th 5, Th 6, Th 7, 8

Proof : Multiplying equation (47) by $u^a u^b$, we get

$$\begin{aligned} u^a u^b (T_{ab;c})h^c + u^a u^b (T_{db})(h^d;_a) + (T_{ad})(h^d;_b) \\ = 2\alpha u^a u^b T_{ab}. \end{aligned} \quad 48$$

The three terms on left hand side of (48) are simplified as follows :

Consider

$$\begin{aligned} u^a u^b (T_{ab;c})h^c &= u^a u^b [(\rho + \mu h^2)u_a u_b \\ &\quad - \frac{1}{2}\mu h^2 g_{ab} - \mu h_a h_b];_c h^c, \\ &= [(\rho + \mu h^2);_c u^a u^b u_a u_b + (\rho + \mu h^2)u_a;_c \\ &\quad u_b u^a u^b + (\rho + \mu h^2)u_b;_c u_a u^a u^b]h^c \\ &\quad - \frac{1}{2}[\mu;_c h^2 g_{ab} u^a u^b + h^2;_c \mu g_{ab} u^a u^b]h^c \\ &\quad - [\mu;_c h_a h_b u^a u^b + \mu h_a;_c h_b u^a u^b + \\ &\quad + \mu h_b;_c u^a u^b]h^c. \end{aligned} \quad 49$$

If we replace $u^a u_a = 1$, $u^a h_a = 0$, $\dot{u}^a u_a = 0$ in (49), we get

$$\begin{aligned} u^a u^b (T_{ab;c})h^c &= (\rho + \mu h^2);_c h^c \\ &\quad - \frac{1}{2}(\mu;_c h^2 h^c + h^2;_c \mu h^c) \\ &= [(\rho + \mu h^2);_c - \frac{1}{2}(\mu h^2);_c]h^c \\ &= [\rho + \frac{1}{2}\mu h^2];_c h^c. \end{aligned} \quad 50$$

Now consider

$$\begin{aligned} u^a u^b (h^d;_a)T_{db} &= u^a u^b h^d;_a [(\rho + \mu h^2)u_d u_b \\ &\quad - \frac{1}{2}\mu h^2 g_{db} - \mu h_d h_b], \end{aligned}$$

$$\begin{aligned}
&= (\rho + \mu h^2) h^d{}_{;a} u^a u^b u_d u_b \\
&\quad - \frac{1}{2} \mu h^2 h^d{}_{;a} u^a u^b g_{db} - \mu h^d{}_{;a} u^a u^b u_d h_b .
\end{aligned} \quad 51$$

Using $u^a u_a = 1$, $u^a h_a = 0$, $h^d{}_{;a} u^a = \dot{h}^d$, in (51), we get

$$\begin{aligned}
u^a u^b (h^d{}_{;a}) T_{db} &= (\rho + \mu h^2) \dot{h}^d u_d - \frac{1}{2} \mu h^2 \dot{h}^d u_d , \\
&= (\rho + \frac{1}{2} \mu h^2) \dot{h}^d u_d .
\end{aligned} \quad 52$$

Further consider,

$$\begin{aligned}
u^a u^b (h^d{}_{;b}) T_{ad} &= u^a u^b (h^d{}_{;b}) [(\rho + \mu h^2) u_a u_d \\
&\quad - \frac{1}{2} \mu h^2 g_{ad} - \mu h_a h_d] \\
&= (\rho + \mu h^2) h^d{}_{;b} u^a u^b u_a u_d \\
&\quad - \frac{1}{2} \mu h^2 h^d{}_{;b} u^a u^b g_{ad} - \mu h^d{}_{;b} u^a u^b h_a h_d ,
\end{aligned} \quad 53$$

We have $u^a u_a = 1$, $u^a h_a = 0$, $h^d{}_{;a} u^a = \dot{h}^d$. Hence from (53), we get

$$\begin{aligned}
u^a u^b (h^d{}_{;b}) T_{ad} &= (\rho + \mu h^2) \dot{h}^d u_d - \frac{1}{2} \mu h^2 \dot{h}^d u_d . \\
&= (\rho + \frac{1}{2} \mu h^2) \dot{h}^d u_d .
\end{aligned} \quad 54$$

Hence from (50), (52), and (54), equation (48) reduces to,

$$\begin{aligned}
&(\rho + \frac{1}{2} \mu h^2)_{;c} h^c + (\rho + \frac{1}{2} \mu h^2) \dot{h}^d u_d \\
&\quad + (\rho + \frac{1}{2} \mu h^2) \dot{h}^d u_d = 2 \alpha u^a u^b T_{ab} .
\end{aligned}$$

$$\begin{aligned}
\text{i.e. } \quad \frac{L}{h} (\rho + \frac{1}{2} \mu h^2) + 2 (\rho + \frac{1}{2} \mu h^2) \dot{h}^d u_d &= 2 \alpha (\rho + \frac{1}{2} \mu h^2) . \\
&\quad \text{vide I (4.2)}
\end{aligned}$$

$$\text{i.e. } \quad \frac{L}{h} e_1 + 2 (\rho + \frac{1}{2} \mu h^2) \dot{h}^d u_d = 2 \alpha e_1 . \quad 55$$

Hence the theorem.

Theorem 6 : If the Magneto-Dust Distribution admits Dynamical Inheritance property along magnetic lines then the necessary and

sufficient condition for the inheritance of active gravitational mass along magnetic lines is given by $\mu (h^2;_c h^c) - \mu h^2 h^a;_a + 2(\rho + \mu h^2) \dot{h}^d u_d = 0$

Proof : On contracting equation (47) with g^{ab} , we get

$$\begin{aligned} (g^{ab} T_{ab};_c) h^c + g^{ab} T_{db} (h^d;_a) + g^{ab} T_{ad} (h^d;_b) \\ = 2 \alpha g^{ab} T_{ab} \end{aligned} \quad 56$$

$$\begin{aligned} \text{i.e. } (\rho;_c) h^c + T^a_d (h^d;_a) + T^b_d (h^d;_b) = 2 \alpha \rho, \quad 57 \\ \text{vide I (4.7)} \end{aligned}$$

$$\text{i.e. } (\rho;_c) h^c + 2 T^{ad} h_{d;a} = 2 \alpha \rho. \quad 58$$

Therefore, by using I (3.4), equation (58) can be written as

$$\begin{aligned} (\rho;_c) h^c + 2 [(\rho + \mu h^2) u^a u^d - \frac{1}{2} \mu h^2 g^{ad} - \mu h^a h^d] h_{d;a} \\ = 2 \alpha \rho. \end{aligned} \quad 59$$

We have $h_{d;a} u^a = \dot{h}^d$, $h^d;_a h^a = -\frac{1}{2} h^2;_a$. Hence from (59), we get

$$\begin{aligned} (\rho;_c) h^c + 2 [(\rho + \mu h^2) h_{d;a} u^a u^d - \frac{1}{2} \mu h^2 h^d;_a g^{ad} \\ - \mu h_{d;a} h^a h^d] = 2 \alpha \rho \end{aligned} \quad 60$$

$$\begin{aligned} \text{i.e. } (\rho;_c) h^c + 2 [(\rho + \mu h^2) \dot{h}_d u^d - \frac{1}{2} \mu h^2 h^d;_a \\ - \mu h_{d;a} h^a h^d] = 2 \alpha \rho \end{aligned}$$

$$\begin{aligned} \text{i.e. } (\rho;_c) h^c + 2 [(\rho + \mu h^2) \dot{h}_d u^d - \frac{1}{2} \mu h^2 h^a;_a \\ - \mu h_{d;a} h^a h^d] = 2 \alpha \rho \end{aligned}$$

$$\begin{aligned} \text{i.e. } \frac{L}{h} \rho - 2 \alpha \rho = \mu h^2 h^a;_a + 2 \mu (-\frac{1}{2} h^2;_a) h^a \\ - 2(\rho + \mu h^2) \dot{h}_d u^d, \end{aligned}$$

$$\text{i.e. } \frac{L}{h} \rho - 2 \alpha \rho = \mu h^2 h^a;_a - \mu h^2;_a h^a - 2(\rho + \mu h^2) \dot{h}_d u^d, \quad 61$$

$$\text{i.e. } \frac{L}{h} \rho - 2 \alpha \rho = \mu h^2 h^a;_a - \mu \frac{L}{h} h^2 - 2(\rho + \mu h^2) \dot{h}_d u^d,$$

$$\text{i.e. } \int_{\mathbf{h}} T - 2 \alpha T = \mu h^2 h^a{}_{;a} - \mu \int_{\mathbf{h}} h^2 - 2 (\rho + \mu h^2) h_d u^d,$$

Hence

$$\int_{\mathbf{h}} T = 2 \alpha T \Leftrightarrow \mu h^2 h^a{}_{;a} - \mu \int_{\mathbf{h}} h^2 - 2 (\rho + \mu h^2) h_d u^d = 0.$$

Hence the proof is complete.

Theorem 7 : If the Magneto-Dust Distribution admits Dynamical Inheritance property along magnetic lines then the following results are equivalent.

$$\text{I) } \int_{\mathbf{h}} \mu = 2 \alpha \mu \quad \text{II) } \int_{\mathbf{h}} h^2 = 0.$$

Proof : If transvect equation (47) by $h^a h^b$, we get

$$\begin{aligned} h^a h^b (T_{ab;c}) h^c + h^a h^b T_{db} (h^d{}_{;a}) + T_{ad} h^a h^b (h^d{}_{;b}) \\ = 2 \alpha h^a h^b T_{ab}. \end{aligned} \quad 63$$

On simplifying the terms on left hand side of (63), we have

Consider

$$\begin{aligned} h^a h^b (T_{ab;c}) h^c &= h^a h^b [(\rho + \mu h^2) u_a u_b \\ &\quad - \frac{1}{2} \mu h^2 g_{ab} - \mu h_a h_b]_{;c} h^c, \\ &\quad \text{vide I (3.4)} \\ &= [(\rho + \mu h^2)_{;c} u_a u_b h^a h^b + (\rho + \mu h^2) u_{a;c} \\ &\quad u_b h^a h^b + (\rho + \mu h^2) u_{b;c} u_a h^a h^b] h^c \\ &\quad - \frac{1}{2} [\mu_{;c} h^2 g_{ab} h^a h^b + \mu h^2_{;c} g_{ab} h^a h^b] h^c \\ &\quad - [\mu_{;c} h_a h_b h^a h^b + \mu h_{a;c} h_b h^a h^b + \\ &\quad + \mu h_{b;c} h_a h^a h^b] h^c. \end{aligned} \quad 64$$

If we substitute $u^a h_a = 0$, $h^a h_a = -h^2$, $u^a u_a = 1$, (64) reduces to

$$\begin{aligned}
h^a h^b (T_{ab;c}) h^c &= -\frac{1}{2} [\mu_{;c} h^2 (-h^2) + \mu h^2_{;c} (-h^2)] h^c \\
&\quad - [\mu_{;c} h^4 - \mu h^2 h_{a;c} h^a \\
&\quad - \mu h^2 h_{b;c} h^b] h^c,
\end{aligned} \tag{65}$$

$$\begin{aligned}
h^a h^b (T_{ab;c}) h^c &= -\frac{1}{2} [\mu_{;c} h^4 - \mu h^2 h^2_{;c}] h^c \\
&\quad (-\mu_{;c} h^4 + \mu h^2 h_{a;c} h^a + \\
&\quad + \mu h^2 h_{b;c} h^b) h^c, \\
&= (\frac{1}{2} \mu_{;c} h^4 + \frac{1}{2} \mu h^2 h^2_{;c}) h^c + \\
&\quad + (-\mu_{;c} h^4 + 2 \mu h^2 h_{a;c} h^a) h^c, \\
&= \frac{1}{2} h^2 [\mu_{;c} h^2 + \mu h^2_{;c}] h^c - \\
&\quad - h^2 \mu_{;c} h^2 h^c - 2 \mu h^2 h_{a;c} h^a h^c, \\
&= \frac{1}{2} h^2 [(\mu h^2)_{;c} h^c] - [\mu_{;c} h^c] h^4 \\
&\quad - 2 \mu h^2 [h_{a;c} h^a] h^c, \\
&= \frac{1}{2} h^2 \left[\frac{L}{h} \mu h^2 \right] - h^2 \left[\frac{L}{h} \mu \right] h^2 \\
&\quad - 2 \mu h^2 [h_{a;c} h^a h^c], \\
&= \frac{1}{2} h^2 \left[\frac{L}{h} \mu h^2 \right] - h^2 \left[\frac{L}{h} \mu \right] h^2 \\
&\quad - 2 \mu h^2 [-\frac{1}{2} h^2_{;c}] h^c, \\
&= (h^2/2) \left\{ \frac{L}{h} \mu h^2 - 2 h^2 \frac{L}{h} \mu + 2 \mu \frac{L}{h} h^2 \right\}, \\
&= (h^2/2) \left\{ \mu \frac{L}{h} h^2 + h^2 \frac{L}{h} \mu - 2 h^2 \frac{L}{h} \mu, \right. \\
&\quad \left. + 2 \mu \frac{L}{h} h^2 \right\} \\
&= (h^2/2) \left[3 \mu \frac{L}{h} h^2 - h^2 \frac{L}{h} \mu \right],
\end{aligned} \tag{66}$$

Now consider

$$\begin{aligned} h^a h^b (h^d;_a) T_{db} &= h^a h^b [(\rho + \mu h^2) u_d u_b \\ &\quad - \frac{1}{2} \mu h^2 g_{db} - \mu h_d h_b] h^d;_a, \end{aligned} \quad 68$$

$$\begin{aligned} &= (\rho + \mu h^2) h^d;_a h^a h^b u_d u_b \\ &\quad - \frac{1}{2} \mu h^2 g_{db} h^a h^b h^d;_a - \mu h^d;_a h^a h^b h_d h_b. \end{aligned} \quad 69$$

Using $h^d;_a h_d = -\frac{1}{2} h^2;_a$, $h^a h_a = -h^2$, $u^a h_a = 0$, $u^a u_a = 1$ in (69), we get

$$\begin{aligned} h^a h^b (h^d;_a) T_{db} &= -\frac{1}{2} \mu h^2 h^a h_d h^d;_a - \mu (-h^2) h^d;_a h_d h^a, \\ &= -\frac{1}{2} \mu h^2 (-\frac{1}{2} h^2;_a) h^a + \mu h^2 (-\frac{1}{2} h^2;_a) h^a \\ &= \frac{1}{4} \mu h^2 h^2;_a h^a - \frac{1}{2} \mu h^2 h^2;_a h^a \\ &= \frac{1}{4} \mu h^2 [h^2;_a h^a - 2 h^2;_a h^a] \\ &= -\frac{1}{4} \mu h^2 h^2;_a h^a \\ &= -\frac{1}{4} \mu h^2 L_h (h^2). \end{aligned} \quad 70$$

Further consider,

$$\begin{aligned} h^a h^b (T_{ad}) (h^d;_b) &= h^a h^b [(\rho + \mu h^2) u_a u_d \\ &\quad - \frac{1}{2} \mu h^2 g_{ad} - \mu h_a h_d] h^d;_b \end{aligned} \quad 71$$

$$\begin{aligned} &= (\rho + \mu h^2) h^d;_b h^a h^b u_a u_d \\ &\quad - \frac{1}{2} \mu h^2 h^d;_b h^a h^b g_{ad} \\ &\quad - \mu h^d;_b h^a h^b h_a h_d, \end{aligned} \quad 72$$

We have $u^a h_a = 0$, $h^a h_a = -h^2$, $h^d;_a h_d = -\frac{1}{2} h^2;_a$.

Hence from (72), we get

$$\begin{aligned} h^a h^b (T_{ad}) (h^d;_b) &= -\frac{1}{2} \mu h^2 h^d;_b h^b h_d + \mu h^2 h^d;_b h^b h_d \\ &= -\frac{1}{2} \mu h^2 (-\frac{1}{2} h^2;_b) h^b \\ &\quad + \mu h^2 (-\frac{1}{2} h^2;_b) h^b \\ &= \frac{1}{4} \mu h^2 L_h h^2 - \frac{1}{2} \mu h^2 L_h h^2 \\ &= -\frac{1}{4} \mu h^2 L_h (h^2). \end{aligned} \quad 73$$

Substituting (67), (70), and (73), in equation (63), we get

$$\begin{aligned} (h^2/2) [3 \mu L_h h^2 - h^2 L_h \mu] - \frac{1}{2} \mu h^2 L_h h^2 &= \\ &= 2 \alpha h^a h^b T_{ab} \end{aligned} \quad 74$$

$$\begin{aligned} (h^2/2) [3 \mu L_h h^2 - h^2 L_h \mu] - \frac{1}{2} \mu h^2 L_h h^2 &= \\ &= 2 \alpha (-\frac{1}{2} \mu h^4) . \end{aligned}$$

$$\begin{aligned} \frac{1}{2} [3 \mu L_h h^2 - h^2 L_h \mu] - \frac{1}{2} \mu L_h h^2 &= \\ &= -\alpha \mu h^2 . \end{aligned}$$

$$\text{i.e. } 3 \mu L_h h^2 - h^2 L_h \mu = -2 \alpha \mu h^2, \quad 75$$

$$\text{i.e. } 2 \mu L_h h^2 - h^2 L_h \mu = -2 \alpha \mu h^2,$$

$$\text{i.e. } h^2 L_h \mu - 2 \mu L_h h^2 = 2 \alpha \mu h^2, \quad 76$$

$$\text{i.e. } L_h \mu - (2/h^2) \mu L_h h^2 = 2 \alpha \mu . \quad 77$$

This implies that $L_h (\mu) = 2 \alpha \mu$ iff $L_h h^2 = 0$.

Hence the proof.

Theorem 8 : If the Magneto-Dust Distribution admits Dynamical Inheritance property along magnetic lines then the magnetic permeability is invariant along flow lines iff the flow lines are expansion free.

Proof : On transvecting equation (47) by $u^a h^b$, we get

$$\begin{aligned} u^a h^b (T_{ab;c}) h^c + u^a h^b T_{db} (h^d_{;a}) + T_{ad} u^a h^b (h^d_{;b}) \\ = 0 \end{aligned} \quad \text{vide I (4.5)} \quad 78$$

The three terms on left hand side of (78) are simplified as follows :

Consider,

$$\begin{aligned}
 u^a h^b (T_{ab;c}) h^c &= u^a h^b [(\rho + \mu h^2) u_a u_b \\
 &\quad - \frac{1}{2} \mu h^2 g_{ab} - \mu h_a h_b] ;_c h^c, \\
 &= [(\rho + \mu h^2) ;_c u_a u_b u^a h^b + (\rho + \mu h^2) u_a ;_c \\
 &\quad u_b u^a h^b + (\rho + \mu h^2) u_b ;_c u_a u^a h^b] h^c \\
 &\quad - \frac{1}{2} [\mu ;_c h^2 g_{ab} u^a h^b + h^2 ;_c \mu g_{ab} u^a h^b] h^c \\
 &\quad - [\mu ;_c h_a h_b u^a h^b + \mu h_a ;_c h_b u^a h^b + \\
 &\quad + \mu h_b ;_c u^a h^b h^a] h^c.
 \end{aligned} \tag{79}$$

$$\begin{aligned}
 \therefore u^a h^b (T_{ab;c}) h^c &= (\rho + \mu h^2) u_b ;_c h^b h^c + \mu h^2 h_a ;_c h^a h^c \\
 &\quad \because u^a u_a = 1, h^a h_a = -h^2, h_a ;_c u^a h^c = -u_b ;_c h^b h^c \\
 &\quad = \rho u_b ;_c h^b h^c + \mu h^2 u_b ;_c h^b h^c \\
 &\quad - \mu h^2 u_b ;_c h^b h^c \\
 u^a h^b (T_{ab;c}) h^c &= \rho u_b ;_c h^b h^c
 \end{aligned} \tag{80}$$

Now consider,

$$\begin{aligned}
 u^a h^b T_{db} (h^d ;_a) &= u^a h^b [(\rho + \mu h^2) u_d u_b - \frac{1}{2} \mu h^2 g_{db} \\
 &\quad - \mu h_d h_b] h^d ;_a, \\
 &= (\rho + \mu h^2) u^a h^b u_a u_b h^d ;_a \\
 &\quad - \frac{1}{2} \mu h^2 u^a h^b g_{db} h^d ;_a - \\
 &\quad - \mu h_d h_b u^a h^b h^d ;_a \\
 &= -\frac{1}{2} \mu h^2 u^a h_d h^d ;_a + \mu h^2 h_d u^a h^d ;_a \\
 &= \frac{1}{2} \mu h^2 h^d ;_a h_d u^a \quad \because h^a h_a = -h^2 \\
 &= \frac{1}{2} \mu h^2 (-\frac{1}{2} h^2 ;_a) u^a, \\
 &= -\frac{1}{4} \mu h^2 (h^2 ;_a) u^a.
 \end{aligned} \tag{81}$$

$$\therefore u^a h^b T_{db} (h^d ;_c) = -\frac{1}{4} \mu h^2 (h^2)^\cdot \quad \because h^2 ;_a u^a = (h^2)^\cdot \tag{82}$$

Further consider,

$$u^a h^b (T_{ad}) (h^d;_b) = u^a h^b [(\rho + \mu h^2) u_a u_d - \frac{1}{2} \mu h^2 g_{ad} - \mu h_a h_d] h^d;_b, \quad 83$$

$$= u^a h^b h^d;_b (\rho + \mu h^2) u_a u_d - \frac{1}{2} \mu h^2 u^a h^b g_{ad} h^d;_b - \mu u^a h^b h_a h_d h^d;_b, \quad 84$$

$$= (\rho + \mu h^2) h^d;_b u_d h^b - \frac{1}{2} \mu h^2 h^b u_d h^d;_b, \quad 85$$

$$= (\rho + \frac{1}{2} \mu h^2) h^d;_b h^b u_d.$$

Hence from (85), (82) and (80), we can write equation (78) as

$$\rho [u_{b;c} h^b h^c] - \frac{1}{4} \mu h^2 (h^2)^\bullet + [\rho + \frac{1}{2} \mu h^2] [h^d;_b h^b u_d], \quad 86$$

$$= 0$$

$$\text{i.e. } -\rho h_{a;c} u^a h^c - \frac{1}{4} \mu h^2 (h^2)^\bullet + \rho h_{d;b} h^b u^d + \frac{1}{2} \mu h^2 h_{d;b} h^b u^d, \quad 87$$

$$= 0$$

$$\text{i.e. } -\frac{1}{4} \mu h^2 (h^2)^\bullet + \frac{1}{2} \mu h^2 [h_{d;b} h^b u^d] = 0,$$

$$\text{i.e. } -\frac{1}{4} \mu h^2 (h^2)^\bullet + \frac{1}{2} \mu h^2 [-u_{d;b} h^b u^d] = 0, \quad 87$$

$$\text{i.e. } -\frac{1}{4} \mu h^2 (h^2)^\bullet - \frac{1}{2} \mu h^2 [-\dot{\mu} h^2 - \mu h^2 \theta - (\mu/2) (h^2)^\bullet] = 0, \quad \text{vide I (8.7)}$$

$$\text{i.e. } -\frac{1}{4} \mu h^2 (h^2)^\bullet + (\mu/2) h^4 + (\dot{\mu}/2) h^4 \theta + (\mu/4) h^2 (h^2)^\bullet = 0,$$

$$\text{i.e. } (\dot{\mu}/2) h^4 + (\mu/2) h^4 \theta = 0, \quad 88$$

$$\text{i.e. } \dot{\mu} + \mu \theta = 0. \quad 89$$

Hence the proof.

3. DYNAMICAL INHERITANCE FOR CONFORMAL KILLING VECTORS :

Subcase I :

Claim 1 : If u is the Conformal Killing vector then Dynamical Inheritance is transferred into Dynamical Collineation.

Proof : Let u be the Conformal Killing vector

Then we have

$$\theta = 0 \quad \text{vide II (2.6)}$$

$$2 \dot{\mu} h^2 + \mu (\dot{h}^2) = 0 \quad \text{vide II (2.15)}$$

$$\dot{\rho} = \frac{1}{2} \dot{\mu} h^2 \quad \text{vide II (2.16)}$$

$$\dot{\rho} = \frac{1}{2} \dot{\mu} h^2$$

Using these results in equation (9), we get

$$\dot{\rho} = 2 \alpha \rho,$$

$$\text{i.e. } L_u \rho = 2 \alpha \rho. \quad 90$$

If $L_u \rho = 2 \alpha \rho$ [vide (90)], equation (20), yields

$$\frac{1}{2} L_u (\mu h^2) = \alpha \mu h^2,$$

$$\text{i.e. } L_u (\mu h^2) = 2 \alpha \mu h^2. \quad 91$$

Using result (91) in (37), we get

$$2 L_u (\mu h^2) = L_h (\mu h^2) - \dot{\mu} h^2, \quad \because \theta = 0$$

$$\text{i.e. } L_u (\mu h^2) = -\dot{\mu} h^2,$$

$$\text{i.e. } L_u (\mu h^2) = -2 \dot{\rho}, \quad \because \dot{\rho} = \frac{1}{2} \dot{\mu} h^2$$

$$\text{i.e. } \frac{L}{u}(\rho) + L(\mu h^2) = 0, \quad 92$$

$$\text{i.e. } 2[2\alpha\rho] + 2\alpha(\mu h^2) = 0, \quad \text{vide (90), (91)}$$

$$\text{i.e. } 2\alpha\rho + \alpha\mu h^2 = 0,$$

$$\text{i.e. } \alpha(2\rho + \mu h^2) = 0, \quad \because 2\rho + \mu h^2 \neq 0$$

$$\text{i.e. } \alpha = 0$$

This means that the Dynamical Inheritance is transformed into
Dynamical Collineation.

$$\text{i.e. } \frac{L}{u}T_{ab} = 0. \quad 93$$

Subcase II :

Claim 2 : If h is the Conformal Killing vector then magnetic field vector becomes Killing vector if $\psi = 0$.

Proof : Let h be the Conformal Killing vector

Then we have

$$h^b{}_{;b} = 4\psi \quad \text{vide II (2.19)}$$

$$h_a u^a = \psi \quad \text{vide II (2.21)}$$

$$\frac{L}{h}h^2 = 2\psi h^2 \quad \text{vide II (2.23)}$$

$$\dot{\mu} + \mu\theta = 0 \quad \text{vide II (2.27)}$$

$$\dot{\rho} + (\rho + \frac{1}{2}\mu h^2)\theta = 0 \quad \text{vide II (2.28)}$$

$$\rho u_b h^b - \frac{1}{2}\mu_{;b} h^b h^2 = 0 \quad \text{vide I (9.16)}$$

$$\psi[2\rho - 3\mu h^2] = 0 \quad \text{vide II (2.30)}$$

We have equation (62), viz.

$$\frac{L}{h}T - 2\alpha T = \mu h^2 h^a{}_{;a} - \mu \frac{L}{h}h^2 - 2(\rho + \mu h^2) h_a u^a, \quad 94$$

If $h^a{}_{;a} = 4\psi$, $\frac{L}{h}h^2 = 2\psi h^2$, $h_a u^a = \psi$, we have from (94)

$$\frac{L}{h}T - 2\alpha T = \mu h^2 4\psi - \mu 2\psi h^2 - 2(\rho + \mu h^2) 2\psi$$

$$\text{i.e. } \mathcal{L}_h T - 2\alpha T = 4\mu h^2 \psi - 2\mu h^2 \psi - 4\rho \psi - 4\mu h^2 \psi,$$

$$\text{i.e. } \mathcal{L}_h T - 2\alpha T = -2\mu h^2 \psi - 4\rho \psi,$$

$$\text{i.e. } \mathcal{L}_h T - 2\alpha T = -2[2\rho \psi] - 4\rho \psi, \quad \because \psi[2\rho - \mu h^2] = 0$$

$$\mathcal{L}_h T - 2\alpha T = -4\rho \psi - 4\rho \psi,$$

$$\mathcal{L}_h T - 2\alpha T = -8\rho \psi. \quad 95$$

If $\psi = 0$ then magnetic field vector becomes Killing vector.

Hence the claim.

Subcase III:

Claim 3 : If $u = h$ both are Conformal Killing vectors then DI ~~is~~ is transformed into Dynamical Collineation.

Proof : Let $u = h = \xi$ be Conformal Killing vectors

Then we have

$$\theta = 0 \quad \text{vide II (2.32).}$$

$$\psi = 0 \quad \text{vide II (2.33).}$$

$$h^b{}_{;b} = 0 \quad \text{vide II (2.38).}$$

$$\mathcal{L}_h h^2 = 0 \quad \text{vide II (2.39).} \quad 96$$

$$\dot{\mu} = 0 \quad \text{vide II (2.42).}$$

$$\dot{\rho} = 0 \quad \text{vide II (2.40).}$$

$$\mu{}_{;b} h^b = 0 \quad \text{vide II (2.41).}$$

$$\mu (h^2)^\cdot = 0 \quad \text{vide II (2.43).}$$

If we use these results (96), in equations (9), (22) and (37),

we get

$$2 \alpha \rho = 0, \quad 2 \alpha \rho + \alpha \mu h^2 = 0, \quad 2 \alpha \mu h^2 = 0. \quad 97$$

This implies that $\alpha = 0$. 98

Again, if we use results (96) in (55), (61) and (76), we get

$$2 \alpha \rho + \alpha \mu h^2 = 0, \quad 2 \alpha \rho = 0, \quad 2 \alpha \mu = 0. \quad 99$$

This implies that $\alpha = 0$.

This means that Dynamical Inheritance is transformed into

Dynamical Collineation. vide (98), (99)

$$\text{i.e. } \quad \underset{\xi}{L} \text{Tab} = 0. \quad \text{vide (1)} \quad 100$$

Hence the claim.